

# Large N QFT

大阪大学素粒子論研究室B4

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# 概要・動機

## 概要

Large N QFTとは・・・場の数 $N \rightarrow \infty$ の理論

AdS/CFTに関係

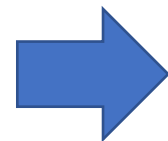
$$\phi^i \quad \phi^{ij} \quad \phi^{ijk}$$

$$i, j, k = 1, 2, \dots, N$$

$$O(N) \quad O(N)^2 \quad O(N)^3 \quad U(N) \quad \dots \text{など}$$

## 動機

- ・大学院でAdS/CFTをやりたい
- ・場の理論の知識がない



Large Nから勉強することに



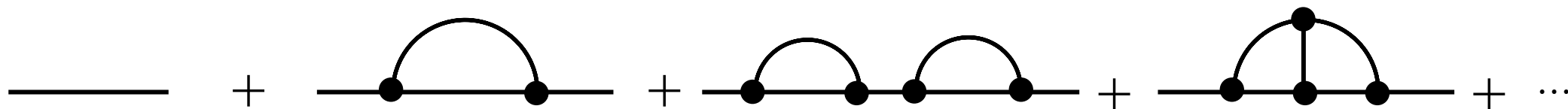
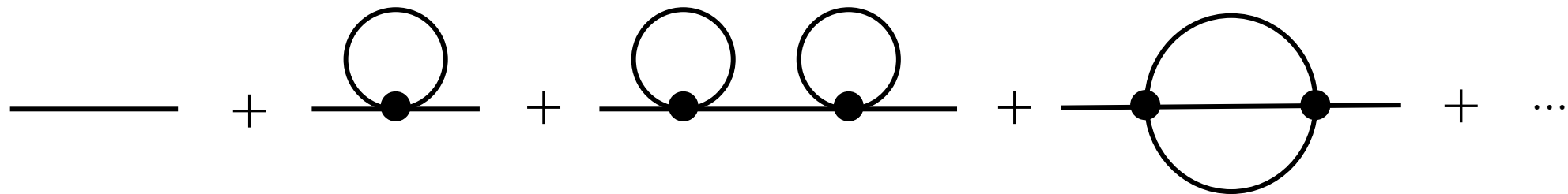
# アウトライン

1. 実スカラー場の理論  $\phi^4 model, \phi^3 model$  (1枚)
2. Large Nアプローチ
  - 2.1 *Vector model* :  $O(N)$   $\phi^i$  (3枚)
  - 2.2 *Matrix model* :  $U(N)$   $\Phi_j^i$  (3枚)
  - 2.2 *Tensor model* :  $O(N)^3$   $\phi^{ijk}$  (3枚)
3. まとめ、今後の展望 (2枚)



# 1 実スカラー場の理論, $\phi^4 model$ , $\phi^3 model$

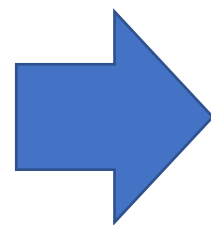
○相互作用が小さいとき→摂動論



○摂動の $n$ 次

→ *Feynman diagram*の数  $\sim n!$

摂動が収束しない



't Hooft

Large N 理論

$\phi^i$     $\Phi_j^i$     $\phi^{ijk}$



## 2.1 Vector model $\phi^i$ ( $O(N)$ )

0+0次元

$$\mathcal{L} = -\frac{1}{2}\phi^i\phi^i - \frac{g}{4!}\phi^i\phi^i\phi^j\phi^j \quad (i, j = 1, 2, \dots, N)$$

$\phi^i$ の添字 $i$ をカラーと呼ぶとする。



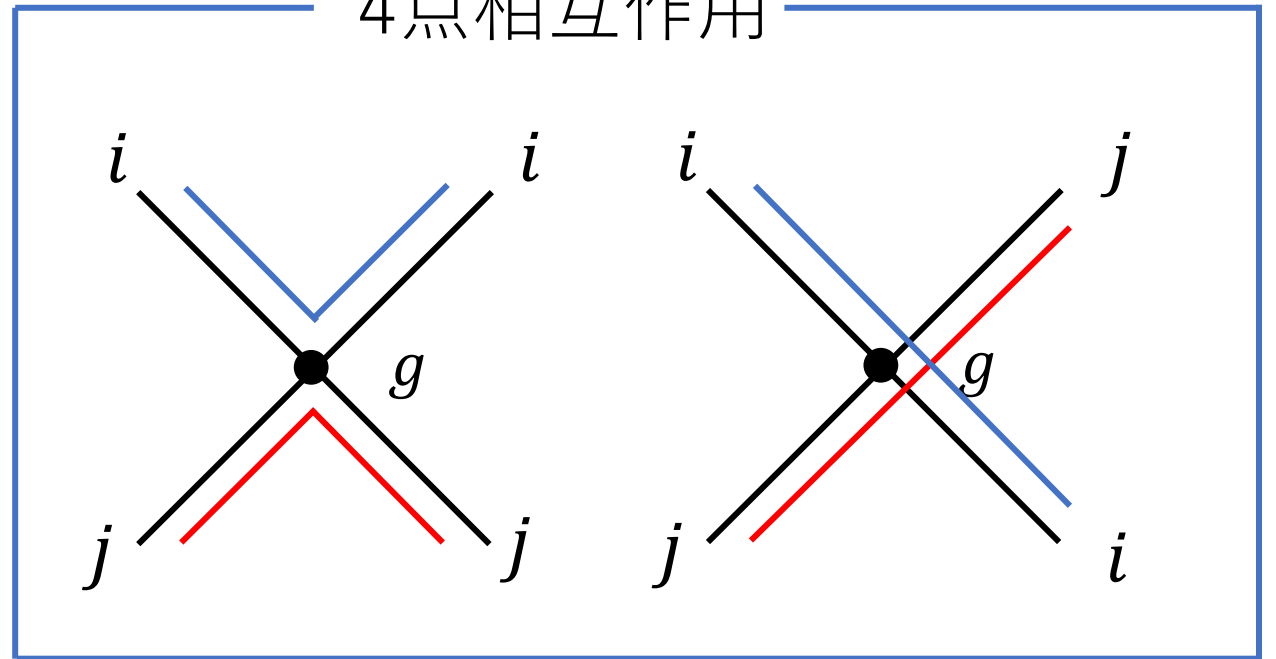
プロパゲーター

$$\langle \phi^i \phi^j \rangle = \delta^{ij}$$



カラーの流れが発生

4点相互作用

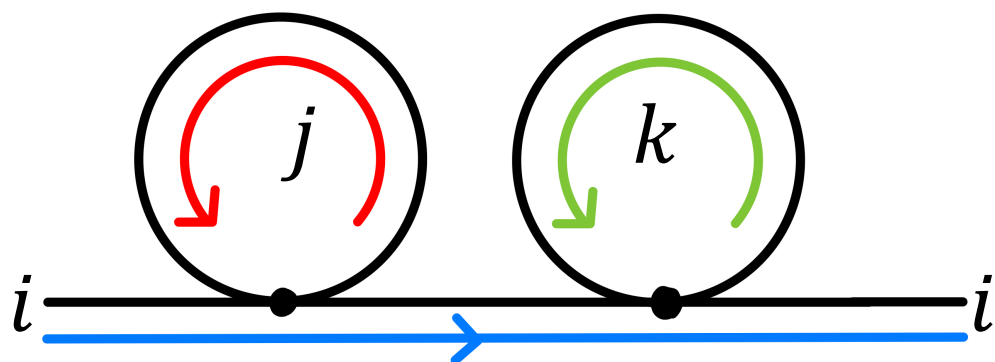




## 2.1 Vector model $\phi^i$ ( $O(N)$ )

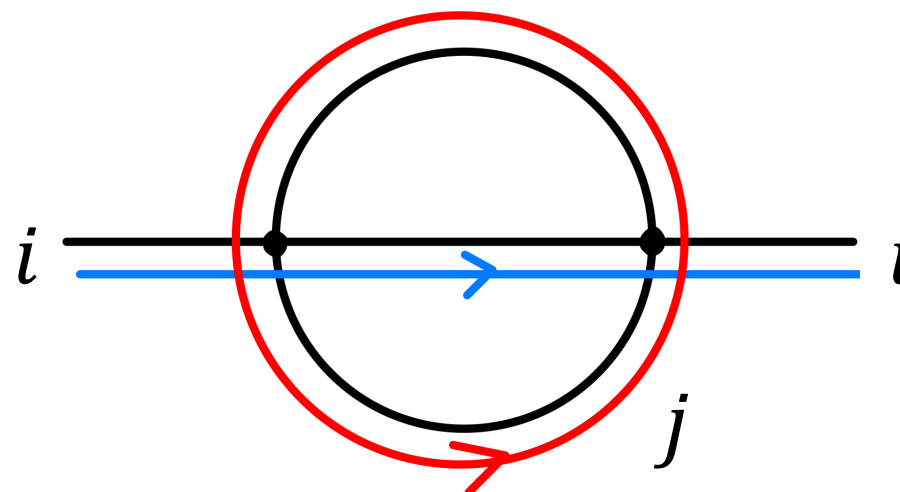
$N \rightarrow \infty$  ,  $\lambda \equiv gN$  fixed (t'Hooft limit)

*snail diagram*



$$\sim g^2 N^2 = \lambda^2$$

*melon diagram*



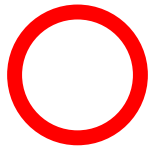
$$\sim g^2 N = \frac{\lambda^2}{N}$$

$$\rightarrow 0 \quad (N \rightarrow \infty)$$

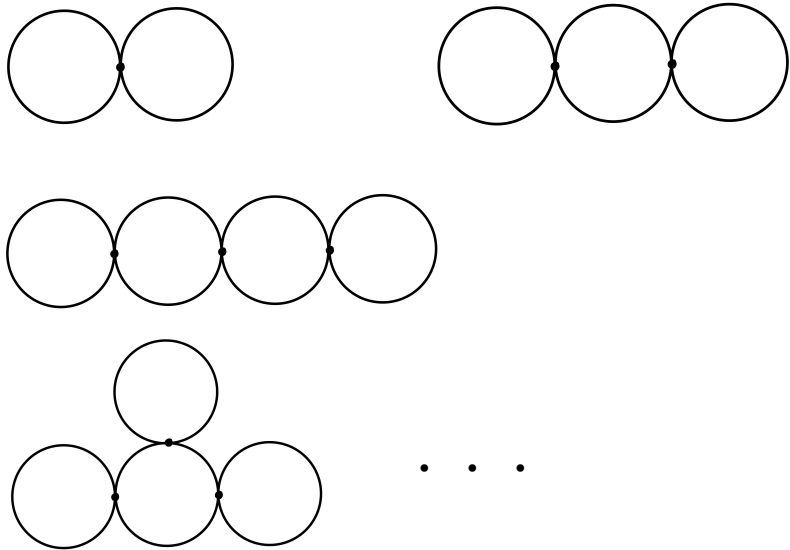


## 2.1 Vector model $\phi^i$ ( $O(N)$ )

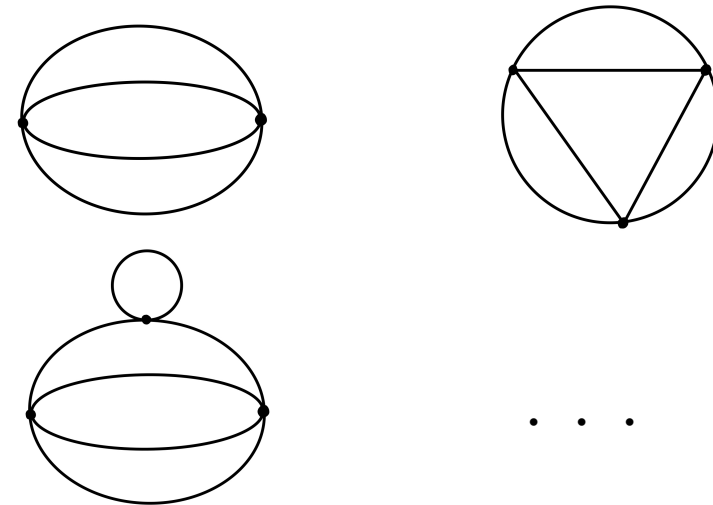
真空ダイアグラムの場合



*bubble diagram*  $\sim O(\lambda^\# N)$



それ以外は高々  $\sim O(\lambda^\#)$



非摂動で解けるモデル

$Z(\vec{J} = 0, \lambda) = \lambda$  の関数



## 2.2 Matrix model $\Phi_j^i$ ( $U(N)$ )

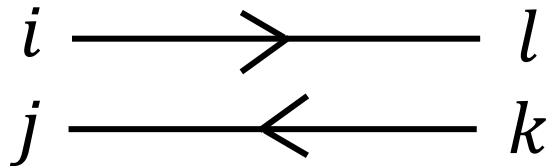
0+0次元

$$\mathcal{L} = -\frac{1}{2} \text{tr}(\Phi^2) - \frac{g}{3!} \text{tr}(\Phi^3) \quad (\Phi: N \times N \text{ matrix})$$

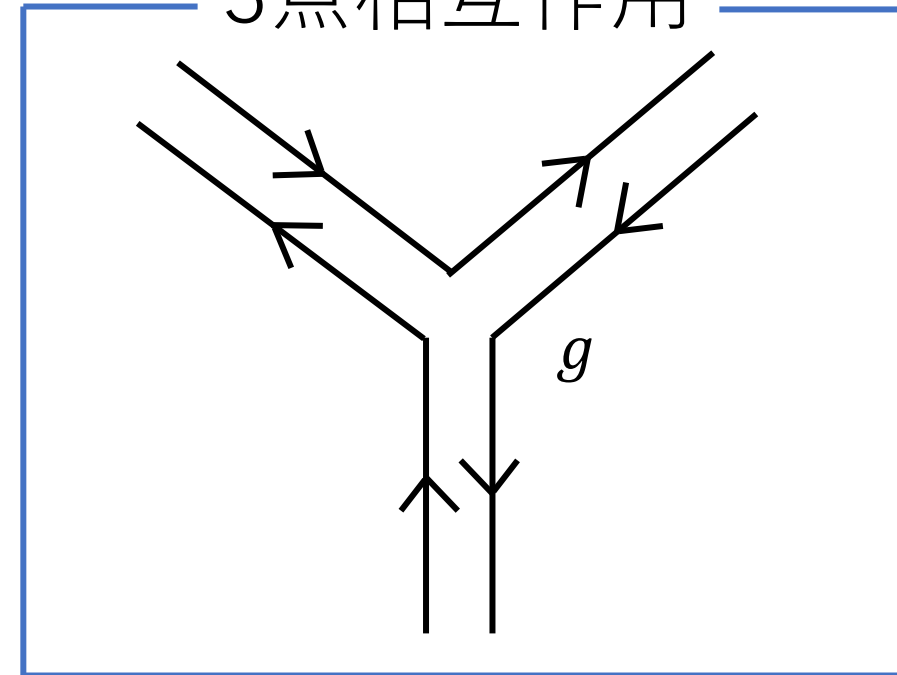


プロパゲーター

$$\langle \Phi_j^i \Phi_l^k \rangle = \delta_l^i \delta_j^k$$



3点相互作用

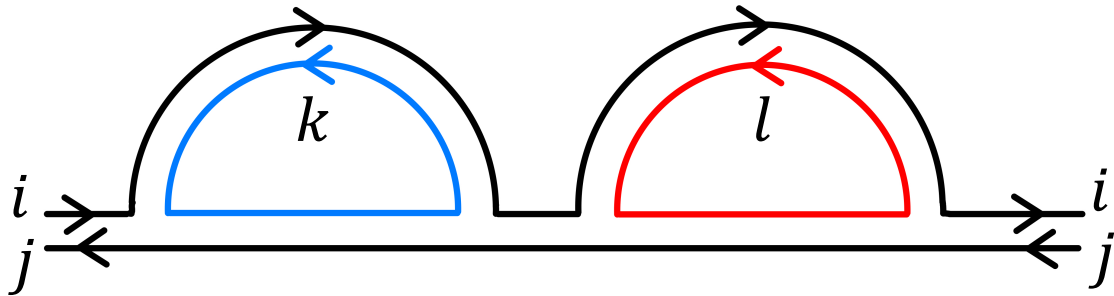




## 2.2 Matrix model $\Phi_j^i (U(N))$

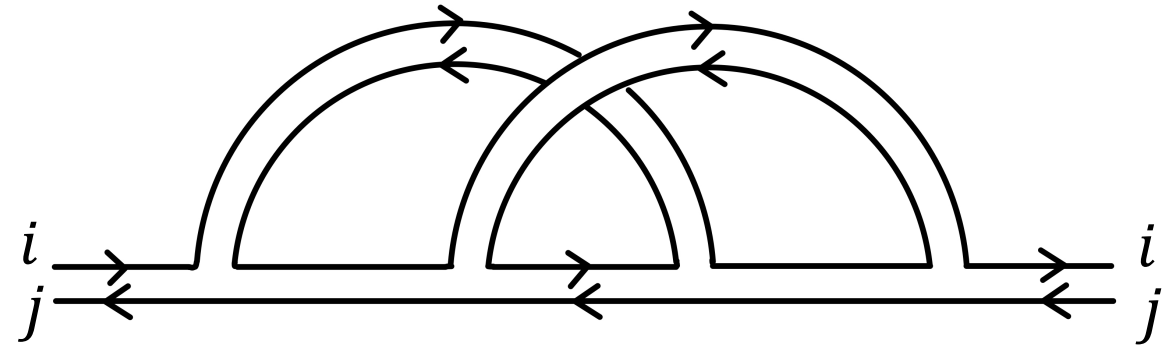
$N \rightarrow \infty$  ,  $\lambda \equiv g^2 N$  fixed (t'Hooft limit)

planar diagram



$$\sim g^4 N^2 = \lambda^2$$

non-planar diagram

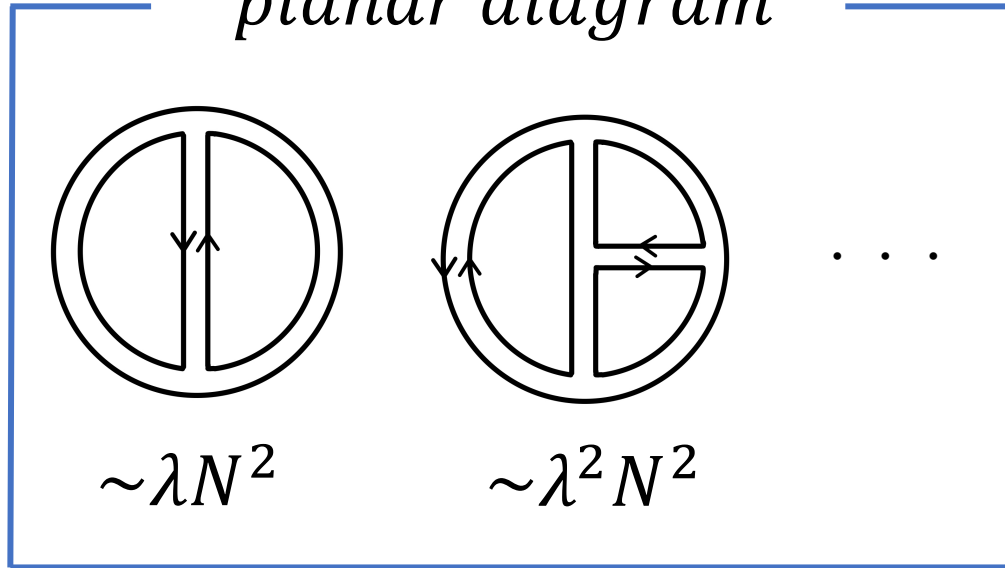


$$\begin{aligned} \sim g^4 &= \lambda^2 N^{-2} \\ &\rightarrow 0 \quad (N \rightarrow \infty) \end{aligned}$$

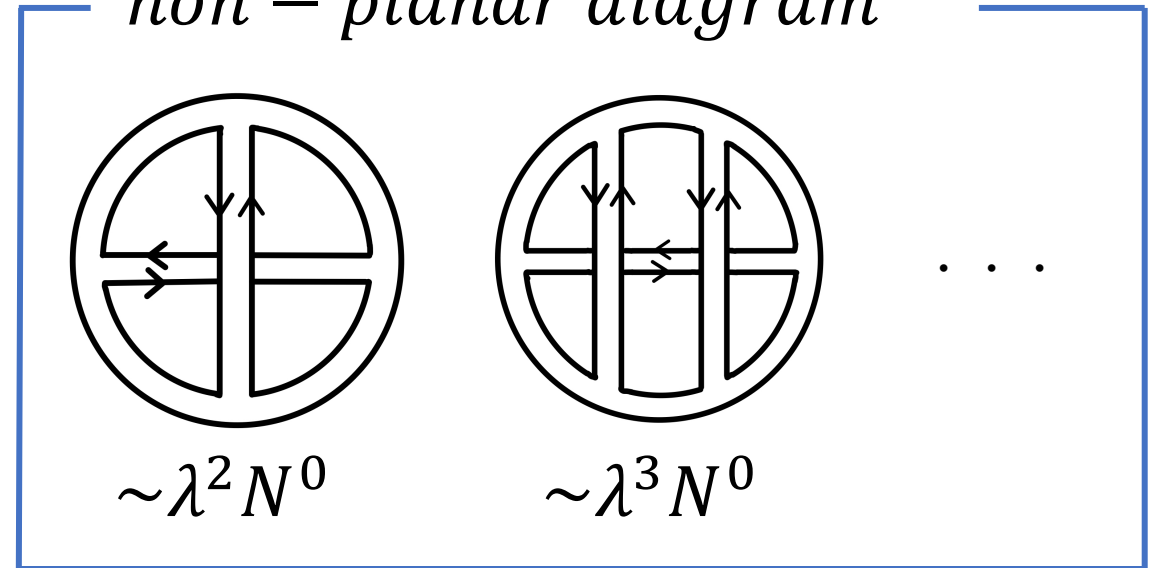


## 2.2 Matrix model $\Phi_j^i (U(N))$

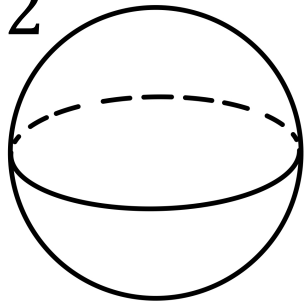
*planar diagram*



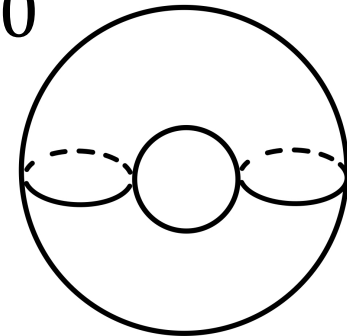
*non – planar diagram*



$\chi = 2$



$\chi = 0$



$\chi = -2$

...

(確率振幅)  $\sim \lambda^\# N^\chi$

( $\chi = 2, 0, -2, \dots$ )

*planar diagram* しか効かない



## 2.3 Tensor model $\phi^{ijk}$ ( $O(N)^3$ )

0+0次元

$$\mathcal{L} = -\frac{1}{2}\phi^{ijk}\phi^{ijk} - \frac{g}{4!}\phi^{i_1j_1k_1}\phi^{i_1j_2k_2}\phi^{i_2j_1k_2}\phi^{i_2j_2k_1}$$

$$(i, j, k, i_1 \dots = 1, 2, \dots N)$$

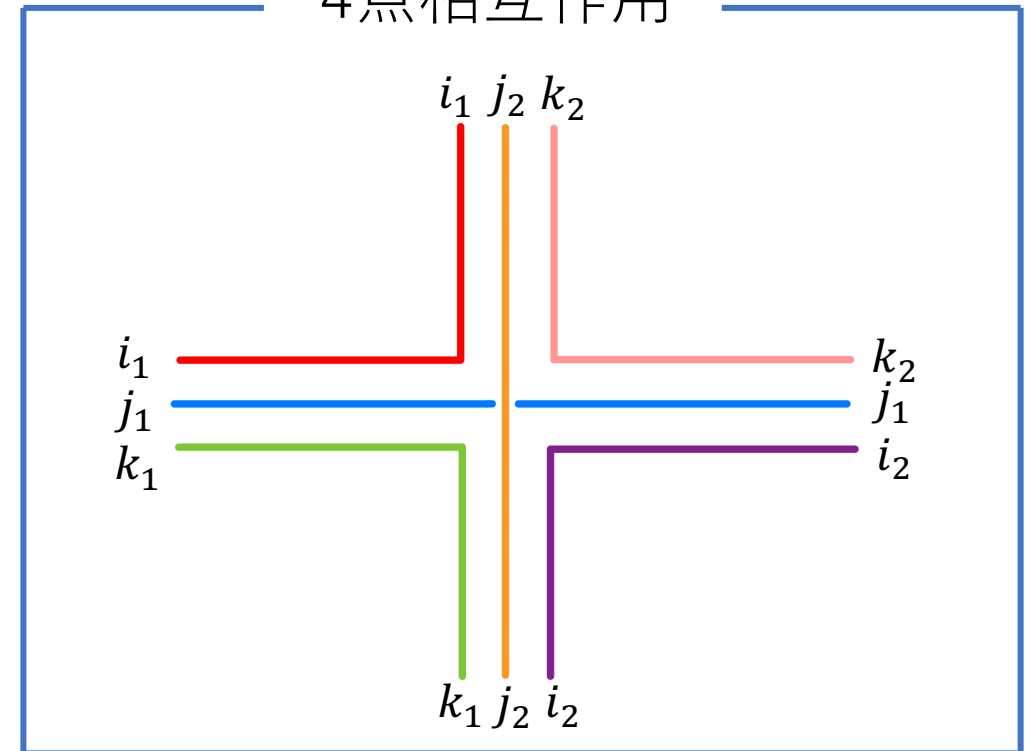


プロパゲーター

$$\langle \phi^{ijk} \phi^{i'j'k'} \rangle = \delta^{ii'} \delta^{jj'} \delta^{kk'}$$



4点相互作用

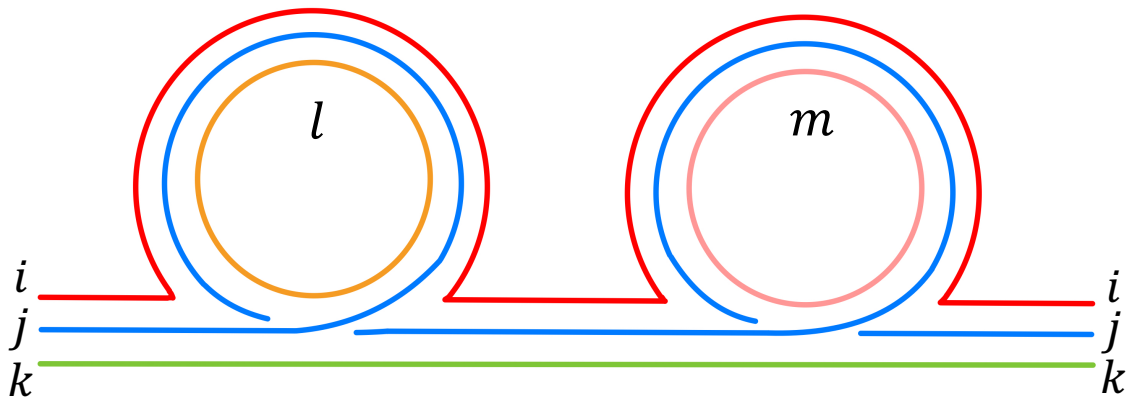




## 2.3 Tensor model $\phi^{ijk}$ ( $O(N)^3$ )

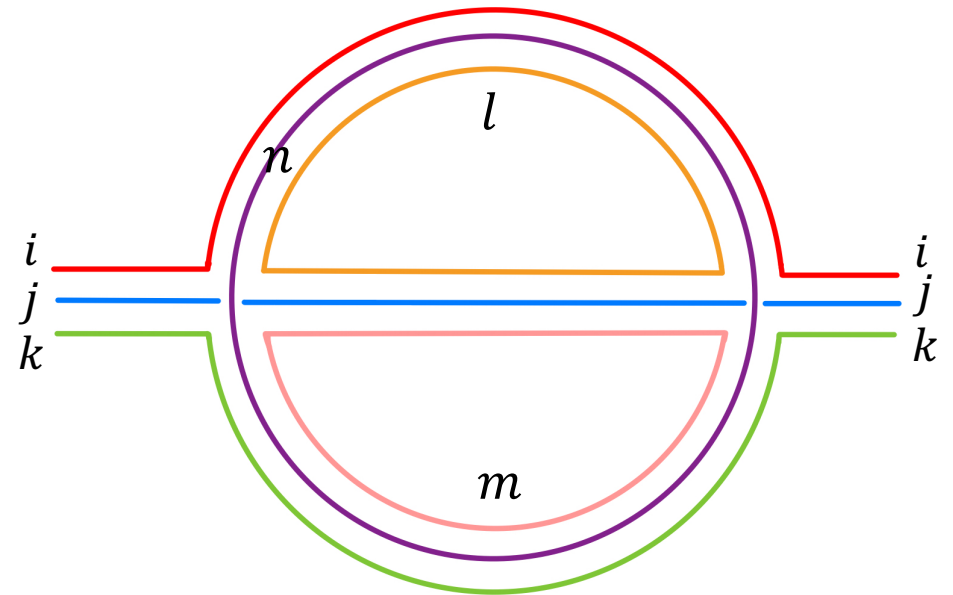
$N \rightarrow \infty$  ,  $\lambda^2 \equiv g^2 N^3$  fixed (t'Hooft limit)

*snail diagram*



$$\sim g^2 N^2 = \frac{\lambda^2}{N}$$
$$\rightarrow 0 \quad (N \rightarrow \infty)$$

*melon diagram*

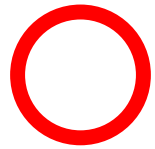


$$\sim g^2 N^3 = \lambda^2$$

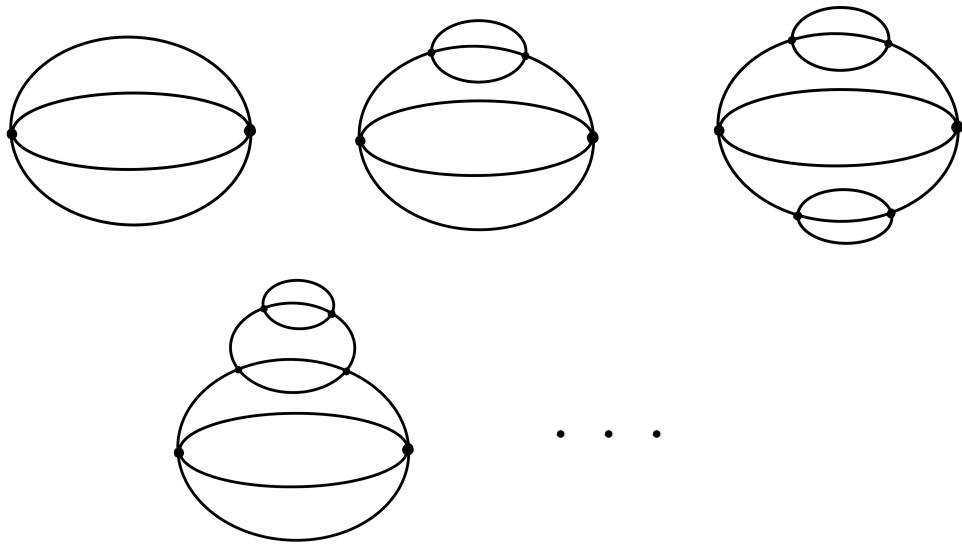


## 2.3 Tensor model $\phi^{ijk}$ ( $O(N)^3$ )

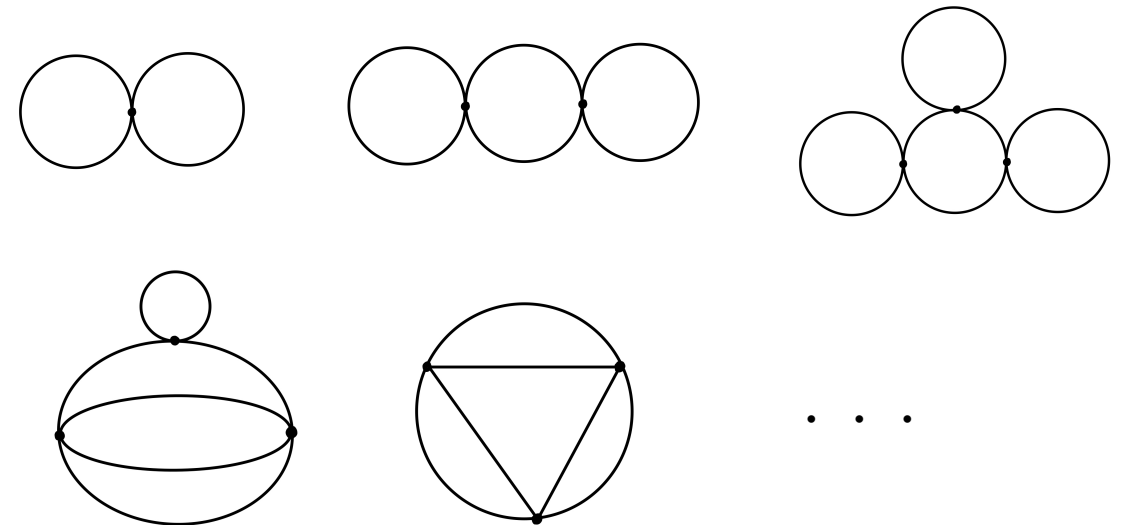
真空ダイアグラムの場合



*melon diagram*  $\sim O(\lambda^\# N^3)$



それ以外は高々  $\sim O(\lambda^\# N^{5/2})$



非摂動で解けるモデル

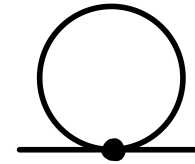
$Z(\vec{J} = 0, \lambda) = \lambda$  の関数



# まとめ

- 場の種類を増やすことでカラーの流れができた。
- 場の数  $N \rightarrow \infty$  の極限では、一部の *Feynman diagram* しか効かない。

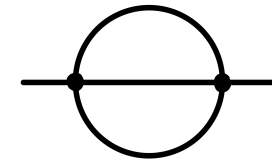
*Vector model*  $\rightarrow$  *snail diagram*



*Matrix model*  $\rightarrow$  *planar diagram*

- トポロジーとの関係が現れた

*Tensor model*  $\rightarrow$  *melon diagram*



- 摂動の収束が良くなり、非摂動で解ける場合もある。

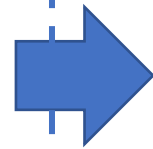


### 3. 今後の展望

卒研で勉強(中)←

→今後勉強 / 研究したい

*Matrix model*

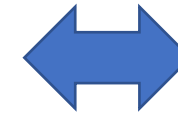


Large N QCD

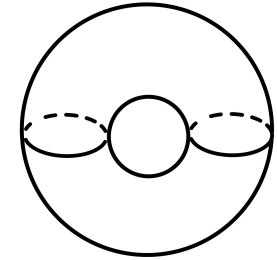
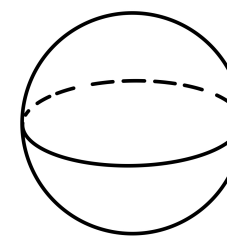
$SU(N)$ ゲージ理論

$$\mathcal{L}_{YM} = -\frac{1}{4} \text{tr}(F_{\mu\nu} F^{\mu\nu})$$

ホログラフィー



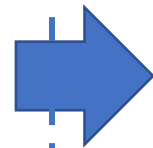
弦理論



...

*Tensor model*

(*melonic*)



Sachdev-Ye-Kitaev

*SYK model*

Large N, melonic,  
solvable

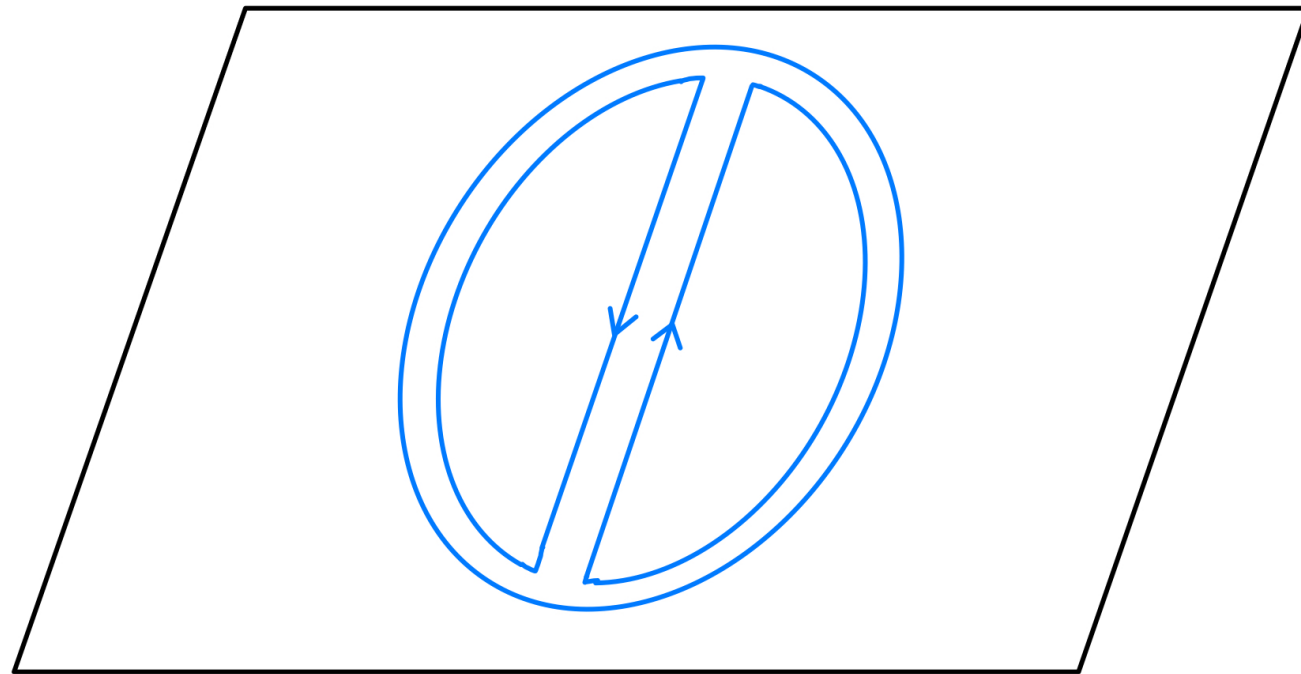
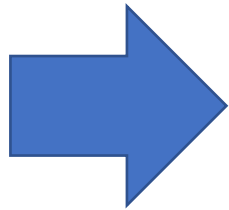
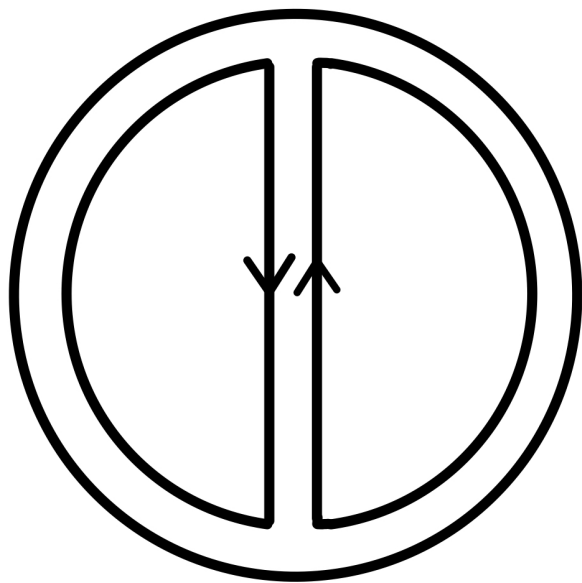


$AdS_2$ 、ブラックホール

2次元重力

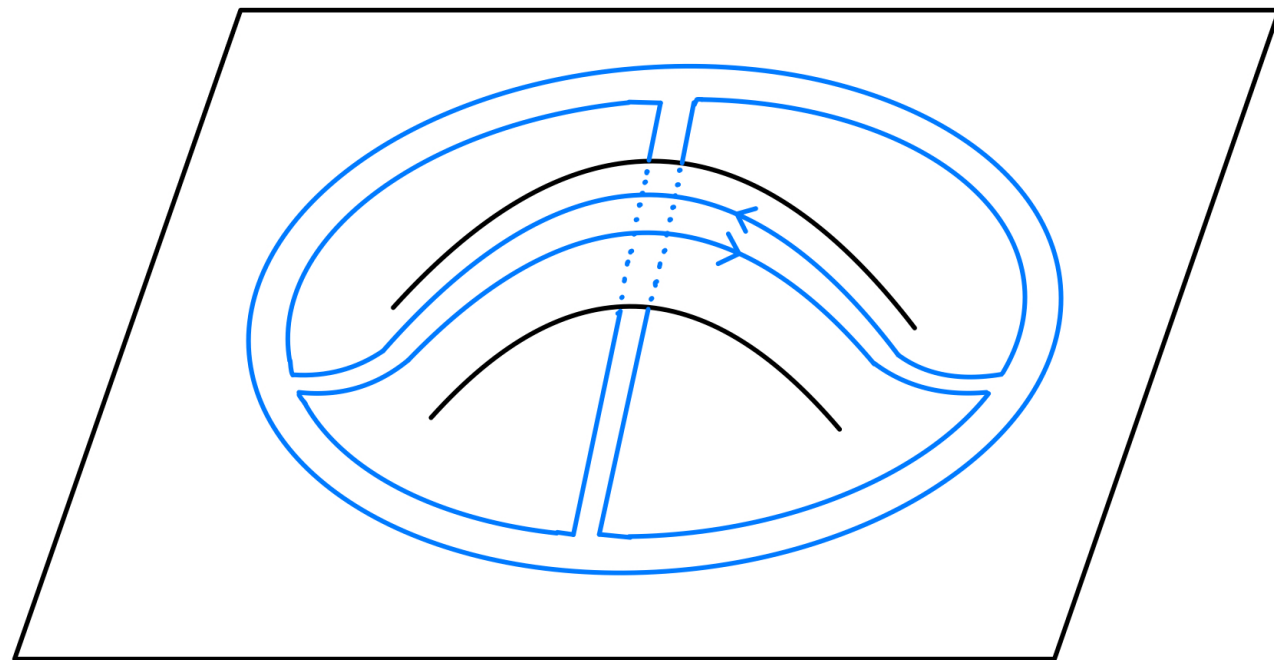
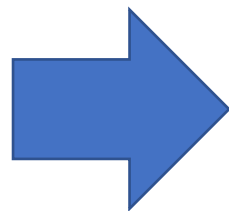
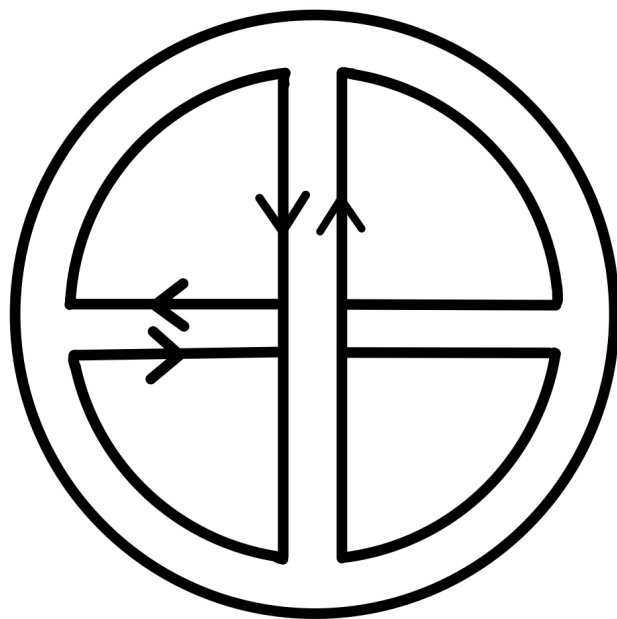


## 4. 補遺





## 4. 補遺






## 4. 補遺

### 2.1 *Vector model* ( $d = 0 + 0$ )

$$\frac{\log Z(\vec{J} = 0)}{N} = \frac{3}{8\lambda} \left( 1 - \sqrt{1 + \frac{2}{3}\lambda} \right)^2 - \frac{1}{2} \log \left( \left( 1 + \sqrt{1 + \frac{2}{3}\lambda} \right) / 2 \right) + o\left(\frac{1}{N}\right)$$



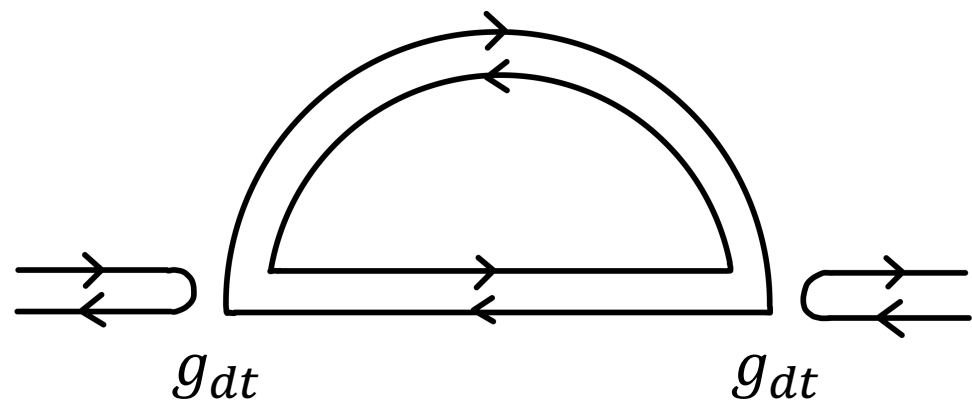
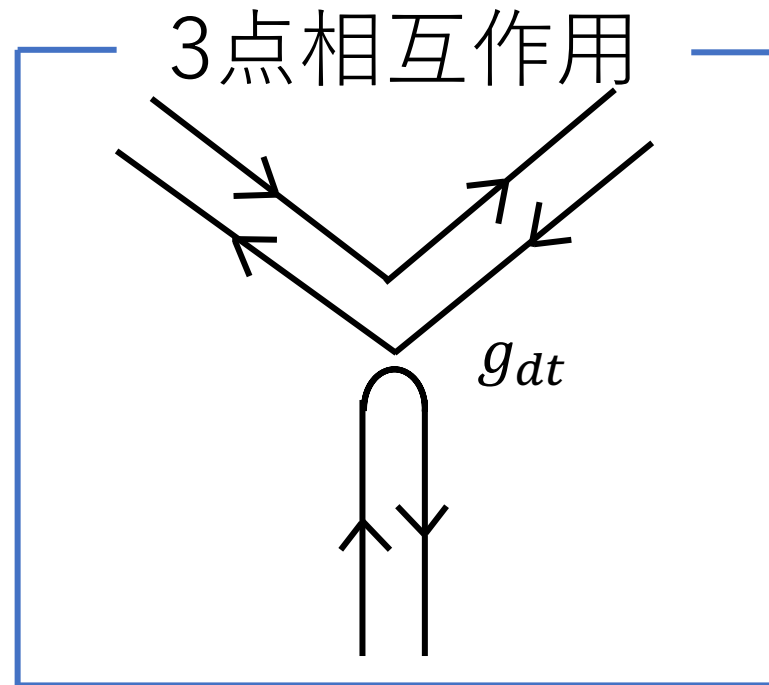
*bubble diagram*



# 4. 補遺

## 2.2 Matrix model

$$\mathcal{L}_{int}^{dt} = -\frac{g_{dt}}{3!} \text{tr}(\Phi^2) \times \text{tr}(\Phi)$$



$$\lambda_{dt} = (g_{dt})^2 N^2$$

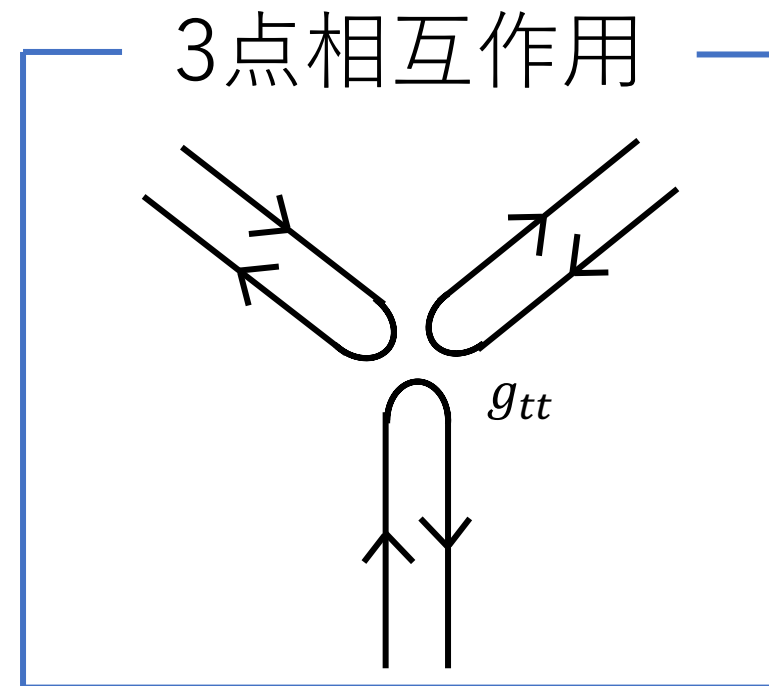
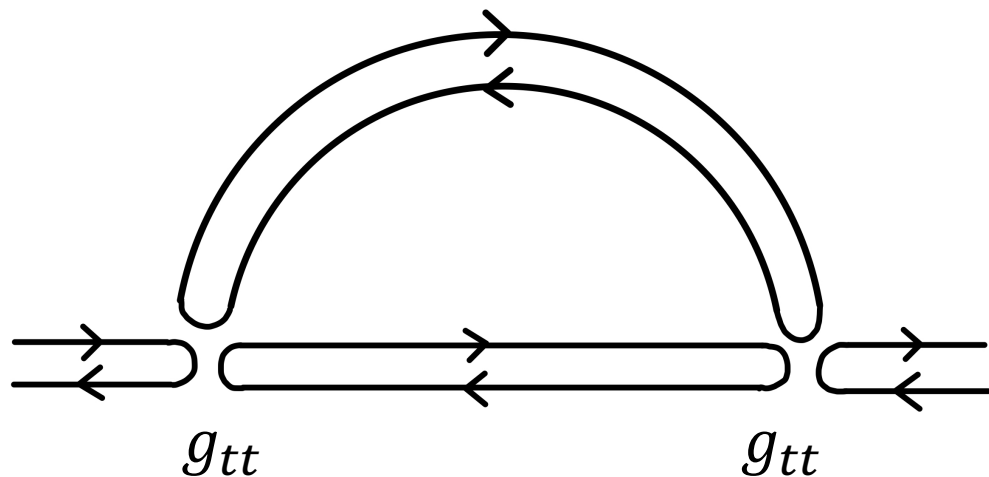
➡  $g_{dt} = \frac{\sqrt{\lambda_{dt}}}{N} \ll g = \frac{\sqrt{\lambda}}{\sqrt{N}}$



# 4. 補遺

## 2.2 Matrix model

$$\mathcal{L}_{int}^{tt} = -\frac{g_{tt}}{3!} \text{tr}(\Phi) \times \text{tr}(\Phi) \times \text{tr}(\Phi)$$



$$\lambda_{tt} = (g_{tt})^2 N^2$$

→

$$g_{tt} = \frac{\sqrt{\lambda_{tt}}}{N} \ll g = \frac{\sqrt{\lambda}}{\sqrt{N}}$$



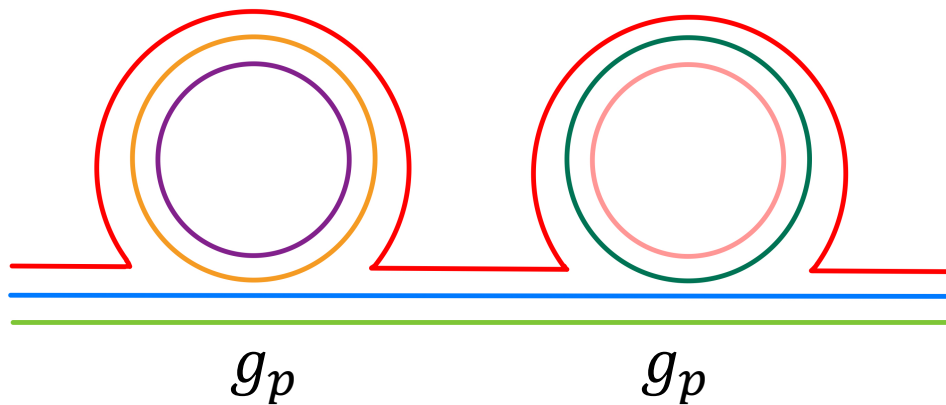
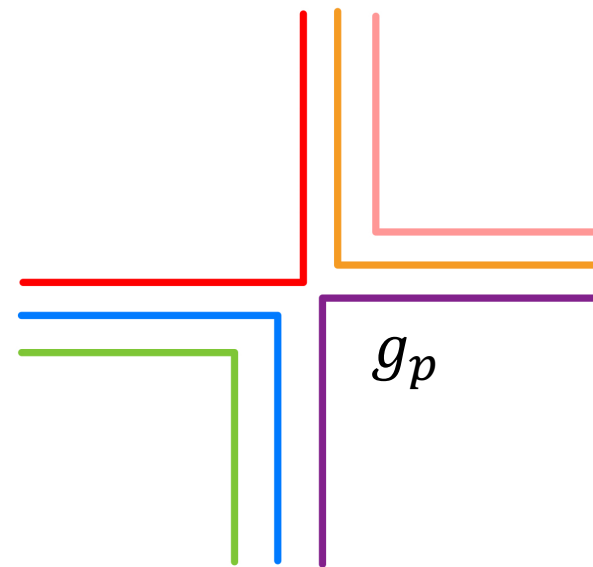
# 4. 補遺

## 2.3 Tensor model

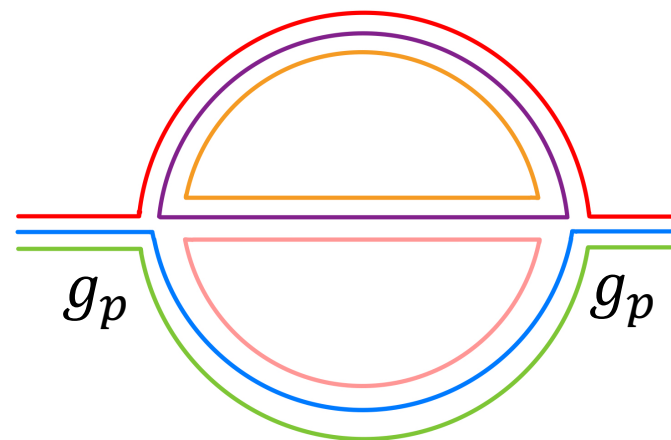
$$\mathcal{L}_{int}^p = -\frac{g_p}{4!} \phi^{i_1 j_1 k_1} \phi^{i_2 j_1 k_1} \phi^{i_2 j_2 k_2} \phi^{i_1 j_2 k_2}$$

$$\lambda_p = g_p N^2$$

$$g_p = \frac{\lambda_p}{N^2} \ll g = \frac{\lambda}{N^{3/2}}$$



$$g_p^2 N^4 = \lambda_p^2$$



$$g_p^2 N^3 = \lambda_p^2 / N$$



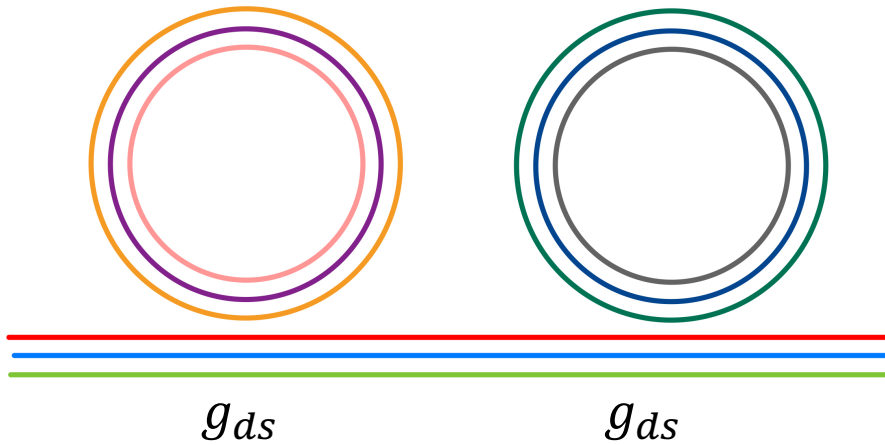
# 4. 補遺

## 2.3 Tensor model

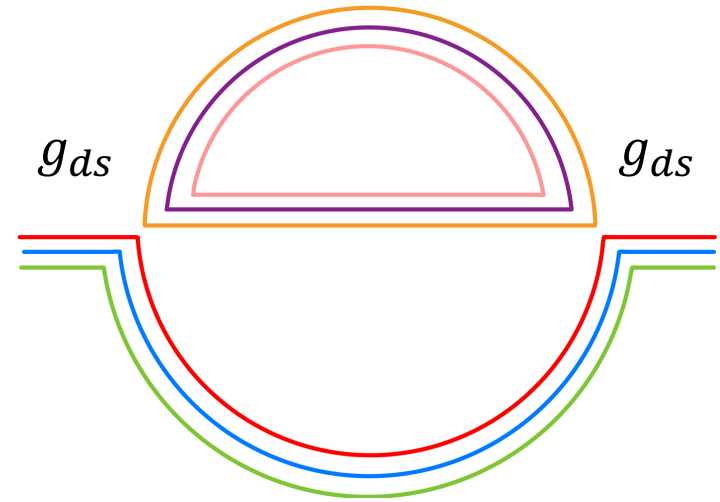
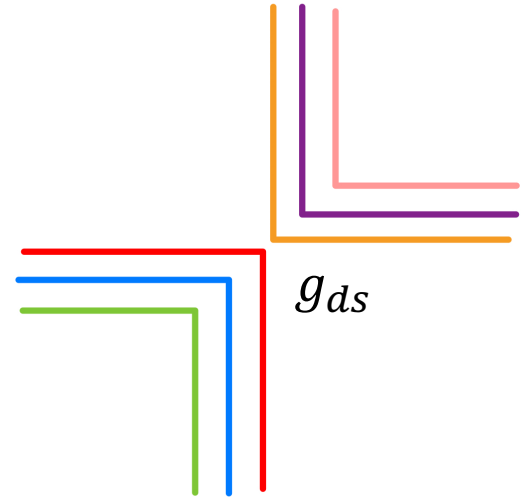
$$\mathcal{L}_{int}^{ds} = -\frac{g^{ds}}{4!} (\phi^{i_1 j_1 k_1} \phi^{i_1 j_1 k_1})^2$$

$$\lambda_{ds} = g_{ds} N^3$$

$$g_{ds} = \frac{\lambda_{ds}}{N^3} \ll g = \frac{\lambda}{N^{3/2}}$$



$$g_{ds}^2 N^6 = \lambda_{ds}^2$$



$$g_{ds}^2 N^3 = \lambda_{ds}^2 / N^3$$